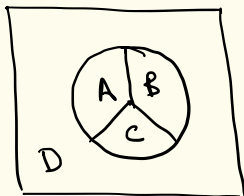


Kitaev-Preskill's Topological Entanglement Entropy



We are given a ground state wave-fn $|\psi\rangle$ of a gapped Hamiltonian.

Consider the combination

$$S_{\text{topo}} = S_A + S_B + S_C - S_{AB} - S_{BC} - S_{CA} + S_{ABC}.$$

AUBUCUD
= whole system.

where S_x denotes vonNeumann entropy of density matrix $\rho_x = \text{Tr}_{\bar{x}} |\psi\rangle\langle\psi|$ and $AB = A \cup B$ etc. All regions A, AB etc. are much larger than the correlation length associated with local operators.

This combination of entropies is 'topological' i.e. it is robust to perturbations of the Hamiltonian (assuming gap doesn't close) or small changes in the geometry of regions A, B, C . e.g. consider deforming region A

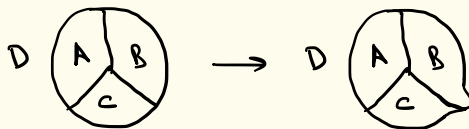
a bit . The change in

$$\begin{aligned} S_{\text{topo}} \text{ is } \Delta S_{\text{topo}} &= \Delta S_A - \Delta S_{AB} - \Delta S_{CA} + \Delta S_{ABC} \\ &= \Delta [S_A - S_{CA}] - \Delta [S_{AB} - S_{ABC}] \end{aligned}$$

$S_A - S_{CA}$ is unchanged since it corresponds to change in entropy associated with attaching C to A , and boundary of C is far from the location where A is

deformed, $\Rightarrow \Delta [S_A - S_{Ac}] \simeq 0$. Similarly
 $\Delta [S_{AB} - S_{ABc}] \simeq 0 \Rightarrow \Delta S_{\text{topo}} \simeq 0$.

Similarly, consider deforming the geometry where three regions meet.



$$\begin{aligned} \text{Now } \Delta S_{\text{topo}} &= \Delta S_B + \Delta S_C - \Delta S_{AB} - \Delta S_{BC} \\ &\quad - \Delta S_{CA} + \Delta S_{ABC} \\ &= \Delta [S_B - S_{AB}] + \Delta [S_{ABC} - S_{BC}] \\ &\quad + \Delta [S_{ABD} - S_{BD}] \end{aligned}$$

where we have used $S_x = S_{\bar{x}}$ for a pure state $\Rightarrow S_C = S_{ABD}$ and $S_{CA} = S_{BD}$.

By the same argument as above, $\Delta S_{\text{topo}} \simeq 0$,
 since all three terms correspond to attaching
 region A to some region and region A is far from
 the location of geometry change.

If S_x for any spatial region x takes the form

$S_x = |\partial x| - \gamma$, we expect γ to be a universal number independent of the shape of region x , as long as x is finite in extent. If so,

$$S_{\text{topo}} = |\partial A| - \gamma + |\partial B| - \gamma + |\partial C| - \gamma \\ - |\partial(AB)| + \gamma - |\partial(BC)| + \gamma - |\partial(CA)| + \gamma$$

$$+ |\partial(CAB)| - \gamma$$

$$= |\partial A| + |\partial B| + |\partial C| - [|\partial A| + |\partial B| - 2|\partial(A \cap B)|]$$

$$- [|\partial B| + |\partial C| - 2|\partial(B \cap C)|] - [|\partial C| + |\partial A| - 2|\partial(C \cap A)|]$$

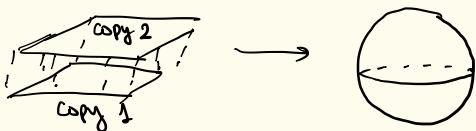
$$+ |\partial A| + |\partial B| + |\partial C| - 2|\partial(A \cap B)| - 2|\partial(B \cap C)|$$

$$- 2|\partial(C \cap A)| - \gamma$$

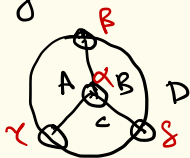
Everything cancels out except $-\gamma$.

$$\Rightarrow \boxed{S_{\text{topo}} = -\gamma}$$

To calculate S_{topo} , we consider two copies of the system and glue them at spatial infinity, thus obtaining a sphere:

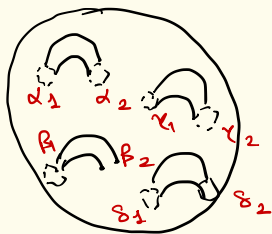


Next we attach wormholes that join the two copies at the intersection of any three regions. There are four such intersections:



The resulting topology is a sphere with four handles:

α_1, α_2 denote the mouth of the wormhole in two copies 1 and 2.



Now,

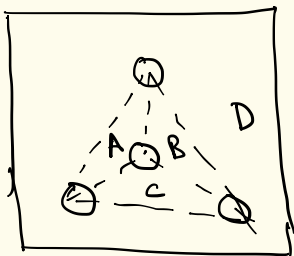
$$2 \text{ Steps} = 4 S_3 - 3 S_4$$

$$\uparrow (S_A + S_B + S_C + S_{ABC})$$

$$\uparrow (S_{AB} + S_{BC} + S_{CA})$$

where $S_A = S_B = S_C = S_{ABC} = S_3$
 = entropy of a region with three punctures

$S_{AB} = S_{BC} = S_{CA} = S_4$
 = entropy of a region with four punctures.

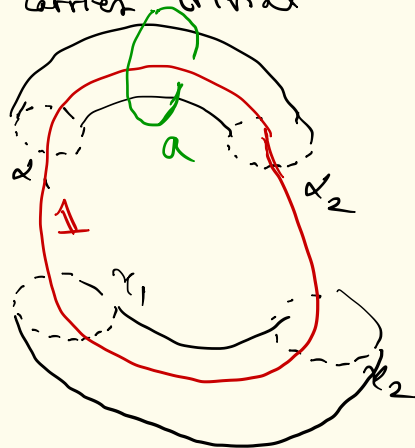


The mouth of each wormhole carries trivial

(i.e. identity) anyon charge. e.g., the loop on the right should detect trivial

anyon charge. This allows

one to calculate the probability that the loop complementary



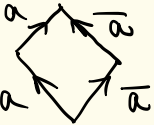
to it, i.e. a loop that has non-zero winding with the above loop carries charge of a specific anyon, say, a . We denote this complimentary loop for the wormhole associated with (α_1, α_2) the **green** loop above.

probability of anyon ' a ' charge at puncture $(\alpha_1, \alpha_2) = |S_{1a}|^2$ where S is the S -matrix. This is because S_{ab} is the amplitude for braiding world-lines of anyon types a and b . Here $b=1$. Using

$$S_{1a} = \frac{da}{D} \quad \text{where } D = \sqrt{\sum_i d_i^2},$$

$$P_a = \frac{d_a^2}{D^2}. \quad \text{Heuristically, we are calculating}$$

the unconditional probability of finding an anyon of type ' a ' at the puncture. This

has the amplitude  =  = da , from our earlier discussion of anyon algebra.

To calculate S_3 , we need to calculate the probability of anyon charges a, b, c at the three punctures. These anyons must fuse to identity, again because the loop that winds around the throat of wormhole must detect $\mathbb{1}$. Denoting this probability as P_{abc} ,

$$P_{abc} = P_a P_b P_{ab \rightarrow \bar{c}} \quad \text{where}$$

$P_{ab \rightarrow \bar{c}}$ is the prob. that a, b fuse to $\bar{c}$ ($=$ prob. that a, b, c fuse to $\mathbb{1}$). To calculate $P_{ab \rightarrow \bar{c}}$, we note,

$$\sum_{ab} P_a P_b P_{ab \rightarrow \bar{c}} = P_c$$

where P_x are the unconditional prob. of finding x . Using $P_a = \frac{d_a^2}{D^2}$ and $d_a d_b = \sum_c N_{ab}^{\bar{c}} d_c$,

one can check that the above eqn is satisfied if

$$\text{we chose, } P_{ab \rightarrow \bar{c}} = \frac{N_{ab}^{\bar{c}} d_c}{d_a d_b}$$

$$\text{Check: } \sum_{ab} \frac{d_a^2 d_b^2}{D^4} \frac{N_{ab}^{\bar{c}} d_c}{d_a d_b} = \sum_b \frac{d_c^2 d_b^2}{D^4} = \frac{d_c^2}{D^2} = P_c$$

$$\begin{aligned}
 \Rightarrow P_{abc} &= P_a P_b P_{ab \rightarrow c} \\
 &= \frac{d_a^2}{D^2} \frac{d_b^2}{D^2} \frac{N_{ab}^c d_c}{d_a d_b} \\
 &= \frac{d_a d_b d_c N_{abc}}{D^4}
 \end{aligned}$$

$$\text{Thus, } S_3 = - \sum_{abc} P_{abc} \log P_{abc}$$

$$= - \sum_{abc} \frac{d_a d_b d_c N_{abc}}{D^4} \log \left[\frac{d_a d_b d_c}{D^4} \right]$$

$$= - \sum_{abc} \frac{d_a d_b d_c N_{abc}}{D^4} [\log(d_a d_b d_c) - 4 \log D]$$

$$= 4 \log D \sum_{ab} \frac{d_a^2 d_b^2}{D^4}$$

$$- 3 \sum_{abc} \frac{d_a d_b d_c N_{abc}}{D^4} \log(d_a)$$

$$= 4 \log D - 3 \sum_{ab} \frac{d_a^2 d_b^2}{D^4} \log(d_a)$$

$$= 4 \log D - 3 \sum P_a \log(d_a)$$

$$\text{where } P_a = \frac{d_a^2}{D^2}.$$

